

## Canonical En

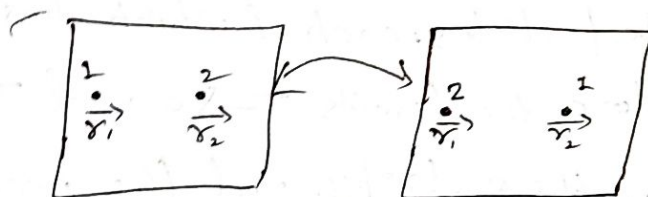
### Gibbs Paradox

Using the concept of microcanonical ensemble, we obtain for an ideal gas (classical)

$$\Sigma(E) = \frac{\pi^{\frac{3N}{2}} V^N (2mE)^{\frac{3N}{2}}}{\Gamma(\frac{3N}{2} + 1)} \quad \text{--- (1)}$$

$\Sigma(E) \rightarrow$  denotes volume in P-space enclosed by a surface of energy E.

Relation (1) leads to several paradoxes if we compute the mixing entropy of identical gases.



Gibbs resolved the paradox in an empirical way. He argued that  $\Sigma(E)$  has been overcounted because permutations of identical particles do not count as an independent state. ~~He then~~ Gibbs assumed  $\Sigma(E)$  is the number of states of the gas with energy less than E,  $N!$  times smaller than ~~the original~~ calculated by the expression (1).

$$\text{Thus } \Sigma(E) = \frac{\pi^{\frac{3N}{2}} V^N (2mE)^{\frac{3N}{2}}}{\Gamma(\frac{3N}{2} + 1) N!} \quad \text{--- (2)}$$

Now the entropy 'S' is given by

$$S = k \ln \Sigma(E) \quad \text{--- (3)}$$

✎ Taking  $\ln N! = N \ln N - N$ , we obtain from (2) and (3)

$$S = Nk \ln \left[ \frac{V}{N} \left( \frac{4\pi m E}{3N} \right)^{3/2} \right] + \frac{5}{2} Nk \quad \text{--- (4)}$$

✎ The modification done by Gibbs does not affect equation of state and other thermodynamic functions of a system.

H.W. ① Using the microcanonical ensemble, compute Helmholtz ~~harmonic~~ free energy  $F(T, N)$  as a function of temperature for a system of  $N$  identical but distinguishable particles, each of which has two energy levels. Explore the limits  $T \rightarrow 0$  and  $T \rightarrow \infty$  of the energy, the entropy and the occupation numbers.

② The internal energy  $E$  of a single component thermodynamic system expressed as a function of entropy  $S$ , and volume  $V$  is of the form

$$E(S, V) = a S^{4/3} V^\alpha, \text{ where } a \text{ and } \alpha \text{ are constants.}$$

Obtain the following:

- (i) value of  $\alpha$ .
- (ii) temperature of the system
- (iii) pressure of the system